

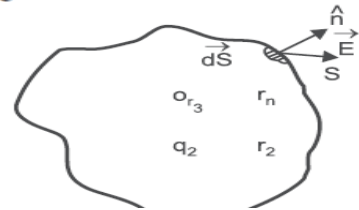
## GAUSS THEOREM

**ELECTRIC FLUX** : The electric flux through a surface is defined as the surface integral of electric field over the surface. It is a measure of the total number of electric field lines passing normal to the surface. It is a scalar. Its SI unit is  $\text{Cm}^2$ .

**Gauss theorem** in electrostatics: It states that the surface integral of electric field over a closed surface is  $\frac{1}{\epsilon_0}$  times the total charge enclosed in the surface.

For surface  $S$

$$\oint_S \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \quad (\text{Total charge inside } S)$$



### Applications of Gauss Theorem

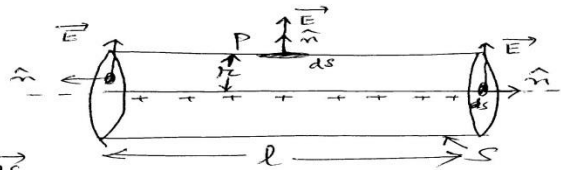
**1. FIELD DUE TO INFINITELY LONG STRAIGHT LINE OF CHARGE:** Consider an infinitely long thin uniformly charged wire. Let P be a point at perpendicular distance  $r$  from the wire.

As the line charge is of infinite expanse, the field lines outside the wire are radial and would be perpendicular to the wire.

Consider a cylindrical gaussian surface  $S$  of radius  $r$  and length  $l$  with its ends perpendicular to the wire as shown below in the fig.

We have;  
Electric flux through the gaussian surface

$$\begin{aligned} \Phi_E &= \oint_S \vec{E} \cdot d\vec{S} \\ &= \int_{\text{circular ends}} \vec{E} \cdot d\vec{S} + \int_{\text{curved surface}} \vec{E} \cdot d\vec{S} \\ &= \int_{\text{circular ends}} E ds \cos 90^\circ + \int_{\text{curved surface}} E ds \cos 0^\circ \\ &= 0 + \int_{\text{curved surface}} E ds = E \int_{\text{curved surface}} ds \\ &= E \cdot 2\pi r l \quad \dots \dots \dots (1) \end{aligned}$$



By Gauss theorem, electric flux through  $S$ ;

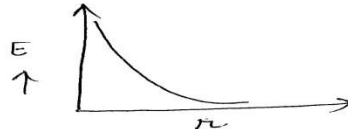
$$\Phi_E = \frac{1}{\epsilon_0} (\text{Charge inside } S) = \frac{1}{\epsilon_0} (\lambda \cdot l) \quad \dots \dots (2)$$

Comparing (1) and (2)

$$E \cdot 2\pi r l = \frac{\lambda l}{\epsilon_0} \Rightarrow E = \frac{\lambda}{2\pi\epsilon_0 r} \Rightarrow E \propto \frac{1}{r}$$

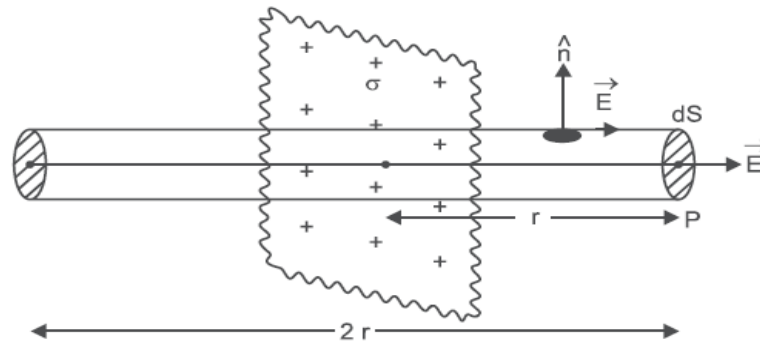
or  $\vec{E} = \frac{\lambda}{2\pi\epsilon_0 r} \hat{r}$

Variation of electric field due to infinite charge with  $\lambda$  and  $r$



## 2. FIELD DUE TO INFINITELY LARGE FLAT SHEET OF CHARGE

**charge:** Consider a large flat sheet of charge density ' $\sigma$ '. Let  $P$  be a point at a distance ' $r$ ' from the sheet.



Take a cylindrical gaussian surface ' $S$ ' with uniform area of cross-section  $dS$  and length  $2r$  symmetrical about the sheet as shown.

Surface integral of the field over the surface  $S$  is given by

$$\begin{aligned}\oint_S \vec{E} \cdot d\vec{S} &= \int_{\text{Curved surface}} \vec{E} \cdot d\vec{S} + \int_{\text{Circular ends}} \vec{E} \cdot d\vec{S} \\ &= \int_{\text{Curved surface}} EdS \cos 90^\circ + \int_{\text{Circular ends}} EdS \cos 0^\circ\end{aligned}$$

**Suppo**

$$\begin{aligned}&= 0 + \int_{\text{Circular ends}} EdS \\ &= E \cdot \int_{\text{Circular ends}} 2dS = E \cdot 2dS \quad \dots(1)\end{aligned}$$

By Gauss theorem

$$\begin{aligned}\oint_S \vec{E} \cdot d\vec{S} &= \frac{1}{\epsilon_0} \quad (\text{charge inside } S) \\ &= \frac{1}{\epsilon_0} (\sigma dS) \quad \dots(2)\end{aligned}$$

From (1) and (2)

$$E \cdot 2dS = \frac{1}{\epsilon_0} (\sigma dS)$$

$$\Rightarrow E = \frac{\sigma}{2\epsilon_0}$$

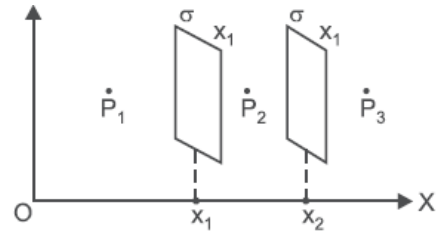
$$\text{or } \vec{E} = \frac{\sigma}{2\epsilon_0} \hat{r} \quad [ \because \text{The field is directed away from the positive charge} ]$$

Suppose there are two large flat sheets of charge with surface charge density  $\sigma$  each as shown, then

Field intensity at  $P_1$  due to the plates  $X_1$  and  $X_2$

$$= \left( \frac{\sigma}{\epsilon_0} + \frac{\sigma}{\epsilon_0} \right) \text{ towards left}$$

$$= -\frac{2\sigma}{\epsilon_0}$$



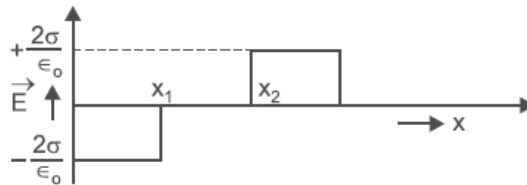
The negative sign indicates that the field is in  $-X$ -direction.

At  $P_2$   $\vec{E} = +\frac{\sigma}{\epsilon_0}$  (due to  $X_1$ )  $-\frac{\sigma}{\epsilon_0}$  (due to  $X_2$ )

$$= \text{Zero.}$$

At  $P_3$   $\vec{E} = +\left| \frac{\sigma}{\epsilon_0} + \frac{\sigma}{\epsilon_0} \right| = \frac{2\sigma}{\epsilon_0}$  in  $+X$ -direction.

The variation of the electric field with  $x$  is represented graphically as under:



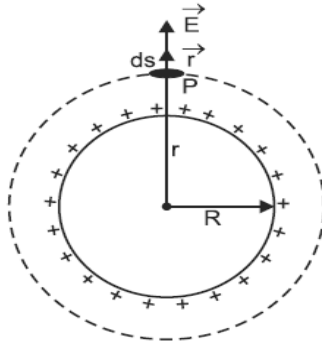
If charge densities are  $\sigma_1$  and  $\sigma_2$  then  $E_1$  (at  $P_1$ ) =  $-(\sigma_1 + \sigma_2)/2\epsilon_0$   
 $E_2 = (\sigma_1 - \sigma_2)/2\epsilon_0$   
 $E_3 = (\sigma_1 + \sigma_2)/2\epsilon_0$

Special case:  $\sigma_1 = \sigma_2 = \sigma$

Then  $E_1 = E_3 = 0$  and  $E_2 = \sigma/\epsilon_0$

### 3. FIELD DUE TO A UNIFORMLY CHARGED SPHERICAL SHELL

Consider a point  $P$  at a distance  $r$  from centre ' $O$ ' of a spherical shell of radius  $R$ . Let  $q$  be the charge on the shell.



Take a spherical Gaussian surface  $S$  with centre  $O$  and radius  $r$ .  
For an area  $ds$  of  $S$ , we have

$$\vec{E} \cdot \vec{ds} = E ds \cos 0^\circ = E ds$$

$$\phi_E = \oint \vec{E} \cdot \vec{ds} = \oint_S E ds = E \cdot 4\pi r^2 \quad \dots(1)$$

By Gauss Theorem,

$$\phi_E = \oint \vec{E} \cdot \vec{ds} = \frac{1}{\epsilon_0} (\text{charge inside } S) \quad \dots(2)$$

From (1) and (2)

$$E \cdot 4\pi r^2 = \frac{1}{\epsilon_0} (\text{charge inside } S)$$

(i) Outside  $\Rightarrow$

$$r > R$$

$$E \cdot 4\pi r^2 = \frac{1}{\epsilon_0} q$$

$$E = \frac{1}{4\pi \epsilon_0} \frac{q}{r^2}.$$

It is same as electric field due to a point charge.

(ii) Inside  $\Rightarrow$

$$r < R$$

$$E \cdot 4\pi r^2 = 0 \quad [ \because q = 0 ]$$

$$E = 0.$$

Electric field inside a charge shell is zero.

